

The Double Copy of a Point Charge

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based on 1912.02177

with Kim, Lee (Daejeon), Nicholson (Edinburgh), Peinador Veiga (QMUL)

and also on previous work with O'Connell (Edinburgh), White (QMUL)

Outline

Part I Brief Introduction to the Double Copy

Part II Detailed Example: Double Copy of Coulomb Solution

Part I

Brief Introduction to the Double Copy

Perturbative gravity is hard!

Feynman rules: expand Einstein-Hilbert Lagrangian $g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}$ [DeWitt '66]

$$\frac{\delta^2 S}{\delta \varphi_{\mu\nu} \delta \varphi_{\sigma'\tau'} \delta \varphi_{\rho''\lambda''}} \rightarrow \text{Sym} \left[-\frac{1}{4} P_3(p \cdot p' \eta^{\mu\nu} \eta^{\sigma\tau} \eta^{\rho\lambda}) - \frac{1}{4} P_6(p^\sigma p^\tau \eta^{\mu\nu} \eta^{\rho\lambda}) + \frac{1}{4} P_3(p \cdot p' \eta^{\mu\sigma} \eta^{\tau\rho} \eta^{\rho\lambda}) + \frac{1}{2} P_6(p \cdot p' \eta^{\mu\nu} \eta^{\sigma\rho} \eta^{\tau\lambda}) + P_3(p^\sigma p^\lambda \eta^{\mu\nu} \eta^{\tau\rho}) - \frac{1}{2} P_3(p^\tau p'^\mu \eta^{\nu\sigma} \eta^{\rho\lambda}) + \frac{1}{2} P_3(p^\rho p'^\lambda \eta^{\mu\sigma} \eta^{\nu\tau}) + \frac{1}{2} P_6(p^\rho p^\lambda \eta^{\mu\sigma} \eta^{\nu\tau}) + P_6(p^\sigma p'^\lambda \eta^{\tau\mu} \eta^{\nu\rho}) + P_3(p^\sigma p'^\mu \eta^{\tau\rho} \eta^{\lambda\nu}) - P_3(p \cdot p' \eta^{\nu\sigma} \eta^{\tau\rho} \eta^{\lambda\mu}) \right],$$

$$\frac{\delta^4 S}{\delta \varphi_{\mu\nu} \delta \varphi_{\sigma'\tau'} \delta \varphi_{\rho''\lambda''} \delta \varphi_{\epsilon'''\kappa'''}} \rightarrow \text{Sym} \left[-\frac{1}{8} P_6(p \cdot p' \eta^{\mu\nu} \eta^{\sigma\tau} \eta^{\rho\lambda} \eta^{\epsilon\kappa}) - \frac{1}{8} P_{12}(p^\sigma p^\tau \eta^{\mu\nu} \eta^{\rho\lambda} \eta^{\epsilon\kappa}) - \frac{1}{4} P_6(p^\sigma p'^\mu \eta^{\nu\tau} \eta^{\rho\lambda} \eta^{\epsilon\kappa}) + \frac{1}{8} P_6(p \cdot p' \eta^{\mu\sigma} \eta^{\nu\tau} \eta^{\rho\lambda} \eta^{\epsilon\kappa}) + \frac{1}{4} P_6(p \cdot p' \eta^{\mu\nu} \eta^{\sigma\tau} \eta^{\rho\lambda} \eta^{\epsilon\kappa}) + \frac{1}{4} P_{12}(p^\sigma p^\tau \eta^{\mu\nu} \eta^{\rho\lambda} \eta^{\epsilon\kappa}) + \frac{1}{2} P_6(p^\sigma p'^\mu \eta^{\nu\tau} \eta^{\rho\lambda} \eta^{\epsilon\kappa}) - \frac{1}{4} P_6(p \cdot p' \eta^{\mu\sigma} \eta^{\nu\tau} \eta^{\rho\lambda} \eta^{\epsilon\kappa}) + \frac{1}{4} P_{24}(p \cdot p' \eta^{\mu\nu} \eta^{\sigma\rho} \eta^{\tau\lambda} \eta^{\epsilon\kappa}) + \frac{1}{4} P_{24}(p^\sigma p^\tau \eta^{\mu\rho} \eta^{\nu\lambda} \eta^{\epsilon\kappa}) + \frac{1}{4} P_{12}(p^\rho p'^\lambda \eta^{\mu\sigma} \eta^{\nu\tau} \eta^{\epsilon\kappa}) + \frac{1}{2} P_{24}(p^\sigma p'^\rho \eta^{\tau\mu} \eta^{\nu\lambda} \eta^{\epsilon\kappa}) - \frac{1}{2} P_{12}(p \cdot p' \eta^{\nu\sigma} \eta^{\tau\rho} \eta^{\lambda\mu} \eta^{\epsilon\kappa}) - \frac{1}{2} P_{12}(p^\sigma p'^\lambda \eta^{\nu\tau} \eta^{\epsilon\mu} \eta^{\rho\kappa}) + \frac{1}{2} P_{12}(p^\sigma p^\rho \eta^{\tau\lambda} \eta^{\mu\nu} \eta^{\epsilon\kappa}) - \frac{1}{2} P_{24}(p \cdot p' \eta^{\mu\nu} \eta^{\tau\rho} \eta^{\lambda\epsilon} \eta^{\kappa\sigma}) - P_{12}(p^\sigma p^\tau \eta^{\nu\rho} \eta^{\lambda\epsilon} \eta^{\kappa\mu}) - P_{12}(p^\rho p'^\lambda \eta^{\nu\tau} \eta^{\epsilon\mu} \eta^{\rho\kappa}) - P_{24}(p^\sigma p'^\rho \eta^{\tau\epsilon} \eta^{\kappa\mu} \eta^{\nu\lambda}) - P_{12}(p^\rho p'^\epsilon \eta^{\lambda\sigma} \eta^{\tau\mu} \eta^{\nu\kappa}) + P_6(p \cdot p' \eta^{\nu\rho} \eta^{\lambda\sigma} \eta^{\tau\epsilon} \eta^{\kappa\mu}) - P_{12}(p^\sigma p^\rho \eta^{\mu\nu} \eta^{\tau\epsilon} \eta^{\kappa\lambda}) - \frac{1}{2} P_{12}(p \cdot p' \eta^{\mu\rho} \eta^{\nu\lambda} \eta^{\sigma\epsilon} \eta^{\tau\kappa}) - P_{12}(p^\sigma p^\rho \eta^{\tau\lambda} \eta^{\mu\epsilon} \eta^{\nu\kappa}) - P_6(p^\rho p'^\epsilon \eta^{\lambda\kappa} \eta^{\mu\sigma} \eta^{\nu\tau}) - P_{24}(p^\sigma p'^\rho \eta^{\tau\mu} \eta^{\nu\epsilon} \eta^{\kappa\lambda}) - P_{12}(p^\sigma p'^\mu \eta^{\tau\rho} \eta^{\lambda\epsilon} \eta^{\kappa\nu}) + 2P_6(p \cdot p' \eta^{\nu\sigma} \eta^{\tau\rho} \eta^{\lambda\epsilon} \eta^{\kappa\mu}) \right].$$

+ infinite number of higher-point vertices. . .



Gravity $\sim$ (Yang-Mills)² in Scattering Amplitudes

Asymptotic states

- Yang-Mills theory: gluon $A_\mu = e^{ik \cdot x} \epsilon_\mu T^a$
colour index a , polarisation ϵ_μ has $D - 2$ dof.

Gravity $\sim (\text{Yang-Mills})^2$ in Scattering Amplitudes

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Contains graviton $h_{\mu\nu}$ + dilaton Φ + B-field $B_{\mu\nu}$, $(D - 2)^2$ dof.

Gravity $\sim$ (Yang-Mills) 2 in Scattering Amplitudes

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Scattering amplitudes

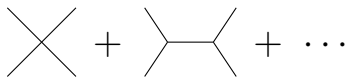
- “Factorisation” of $\epsilon^\mu, \tilde{\epsilon}^\nu$ preserved by interactions!
- **Double copy** $\mathcal{A}_{\text{grav}}(\epsilon_i^{\mu\nu}) \sim (\text{prop})^{-1} \mathcal{A}_{\text{YM}}(\epsilon_i^\mu) \times \mathcal{A}_{\text{YM}}(\tilde{\epsilon}_i^\mu)$ | colour stripped
- Famous application: supergravity UV behaviour. [Bern, Carrasco, Johansson, Roiban, ...]

String theory origin

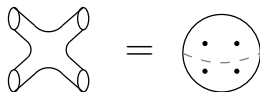
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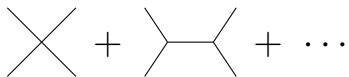
particle scattering
(many Feynman diagrams)



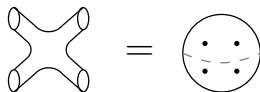
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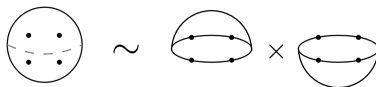
Gravity (closed strings) vs. gauge theory (open strings):

Asymptotic states (vertex operators): $V_{\text{closed}}(\varepsilon^{\mu\nu} = \epsilon^\mu \tilde{\epsilon}^\nu) \sim V_{\text{open}}(\epsilon^\mu) \bar{V}_{\text{open}}(\tilde{\epsilon}^\nu)$

Scattering amplitudes:

KLT relations

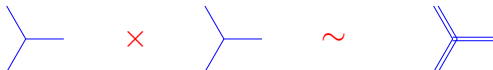
[Kawai, Lewellen, Tye 86]



Field theory limit: **Gravity $\sim$ (Yang-Mills)²** (KLT, BCJ, CHY, ...)

Why simpler?

Basic example: 3-pt interactions.



Gauge theory field A_μ^a

3-pt vertex:
$$f^{abc} V^{\mu\nu\lambda} A_\mu^a(p_1) A_\nu^b(p_2) A_\lambda^c(p_3)$$

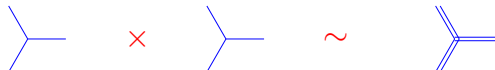
$$V^{\mu\nu\lambda} = (p_1 - p_2)^\lambda \eta^{\mu\nu} + (p_2 - p_3)^\mu \eta^{\nu\lambda} + (p_3 - p_1)^\nu \eta^{\lambda\mu}$$

Gravity field $H_{\mu\mu'} \sim$ graviton + dilaton + B-field 'fat graviton'

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Great simplification: **index factorisation**, c.f. ~ 100 terms in GR 3-pt vertex!

Powerful implementation: **colour-kinematics** duality.

[Bern, Carrasco, Johansson '08] [...]

New directions in (classical) perturbative gravity

Generically, double copy applies in **perturbation theory**.

- **Double-copy-like** field theory for gravity.

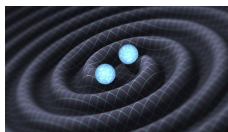
[Goldberger et al] [Luna et al] [Cheung et al] [Plefka et al] [Borsten et al] [...]

- **Gauge-invariant** approach: classical physics from scattering amplitudes.

[Neill et al] [Bjerrum-Bohr et al] [Kosower et al] [Guevara et al] [Huang et al] [Arkani-Hamed et al] [...]

- **Beyond Minkowski**: amplitudes on plane wave backgrounds.

[Adamo et al] [...]

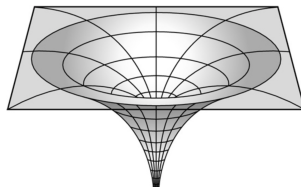


- **Highlight**: new G^3 (3PM) corrections to 2-body potential.

[Bern et al]

Beyond perturbation theory

Question: is a black hole a double copy of something?

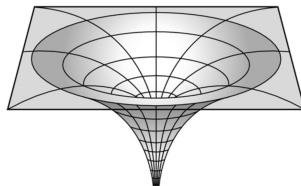


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Challenges:

- What is “graviton” in exact solution?
- Non-perturbative double copy?

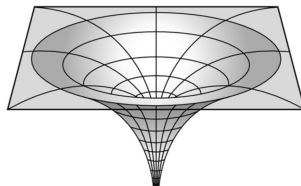


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Still...

- Can relate to perturbation theory.

Examples: Schwarzschild [Duff 73; Neill, Rothstein 13], shockwave [Saitome, Akhoury '12].

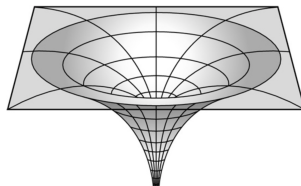
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Stationary Kerr-Schild spacetimes

[RM, O'Connell, White 14]

“Exact perturbation”

$$g_{\mu\nu} = \eta_{\mu\nu} + \phi k_\mu k_\nu$$

where k_μ is null and geodesic wrt $\eta_{\mu\nu}$ and $g_{\mu\nu}$.

$$(k^\mu = g^{\mu\nu} k_\nu = \eta^{\mu\nu} k_\nu)$$

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Einstein equations linearise:

- $g^{\mu\nu} = \eta^{\mu\nu} - \phi k^\mu k^\nu$

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Stationary vacuum case (take $k_0 = 1$): $R^0{}_0 = \frac{1}{2} \nabla^2 \phi = 0$

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“Single copy” linear $\leftrightarrow$ Abelian,

$$A_\mu^a = \phi k_\mu c^a$$

c^a const

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$$0 = D_\mu F^{a\mu\nu} = c^a \begin{cases} -\nabla^2 \phi & \nu = 0 \\ -\partial_\ell [\partial^i (\phi k^\ell) - \partial^\ell (\phi k^i)] & \nu = i \end{cases}$$



Simplest example: point charge

Check spherically symmetric solutions sourced by point charge.

Einstein theory: Schwarzschild solution

- $$g_{\mu\nu} = \eta_{\mu\nu} + \phi k_\mu k_\nu, \quad \phi(r) = \frac{2M}{r}, \quad k = dt + dr$$

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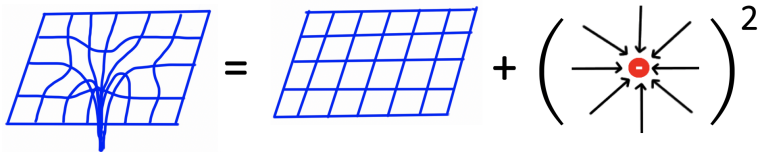
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Schwarzschild $\sim$ (Coulomb)²



Many more examples

Rotation

- Kerr black hole: (M, a) , $a = J/M$.
- $D > 4$: Myers-Perry black holes (M, a_i) .
Other black holes families? Black rings, etc...

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Cosmological constant $\leftrightarrow$ constant charge density. [Luna, RM, O'Connell, White '15]
[Bahjat-Abbas, Luna, White '17; Carrillo-Gonzalez, Penco, Trodden 17]

NUT charge $\leftrightarrow$ magnetic monopole: multi-Kerr-Schild [Luna, RM, O'Connell, White '15]

$$g_{\mu\nu}^{(\text{Taub-NUT})} = \eta_{\mu\nu} + \phi k_{\mu} k_{\nu} + \psi \ell_{\mu} \ell_{\nu}, \quad \phi \propto M, \quad \psi \propto N \Rightarrow A_{\mu}^{(\text{dyon})} = \phi k_{\mu} + \psi \ell_{\mu}$$

Radiation from accelerated particle: correct Bremsstrahlung.

[Luna, RM, Nicholson, O'Connell, White '16]

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Weyl (spinorial) double copy for vacuum type D metrics and pp-waves.
[Luna, RM, Nicholson, O'Connell '18]

Many others – today's talk by Tucker Manton

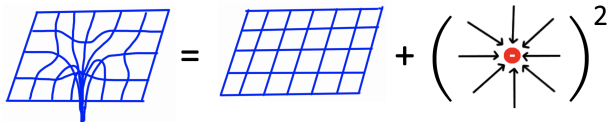
Part II

Double Copy of Coulomb Solution – Revisited

[Kim, Lee, RM, Nicholson, Peinador Veiga '19]

Part II Double Copy of Coulomb Solution

Simplest example: (Coulomb)² $\sim$ Schwarzschild.

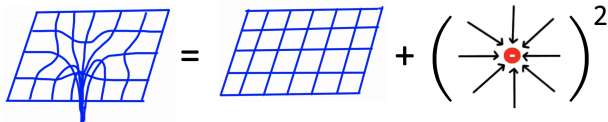


But (YM)² $\sim$ Einstein $h_{\mu\nu}$ + dilaton Φ + B-field $B_{\mu\nu}$.

Other fields?

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But $(\text{YM})^2 \sim \text{Einstein } h_{\mu\nu} + \text{dilaton } \Phi + \text{B-field } B_{\mu\nu}$.

Other fields?

Double-copy structure of Einstein equations?

double copy

Gravity = $\text{YM} \times \text{YM}$



double field theory

doubled geometry $(x^\mu, \tilde{x}_\mu)$

Double copy of Coulomb: perturbative approach

Double copy for Coulomb?

Linearised “fat graviton”:

[Luna, RM, Nicholson, Ochirov, O’Connell, White, Westerberg 16]

$$H_{\mu\nu} = \text{graviton } h_{\mu\nu} + \text{dilaton } \phi + \text{B-field } B_{\mu\nu}$$

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Coulomb usual gauge:

$$A_\mu = -\frac{q^a}{r} u_\mu$$

$$u^\mu = (1, 0, 0, 0).$$

Natural double copy:

[also Goldberger, Ridgway 16]

$$H_{\mu\nu} = \frac{M}{r} u_\mu u_\nu$$

both graviton and dilaton.

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Coulomb different gauge:

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$$k = dt + dr, \quad k^2 = 0.$$

Natural double copy:

[RM, O’Connell, White 14]

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exact Schwarzschild!
Kerr-Schild double copy.

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$$A_\mu = -\frac{q^a}{r} u_\mu$$

$$u^\mu = (1, 0, 0, 0).$$

Natural double copy:

[also Goldberger, Ridgway 16]

$$H_{\mu\nu} = \frac{M}{r} u_\mu u_\nu$$

both graviton and dilaton.

Coulomb different gauge:

$$A_\mu = \frac{q}{r} k_\mu$$

$$k = dt + dr, \quad k^2 = 0.$$

Natural double copy:

[RM, O’Connell, White 14]

$$h_{\mu\nu} = \frac{2M}{r} k_\mu k_\nu$$

exact Schwarzschild!
Kerr-Schild double copy.

Both consistent with other double copy approaches,
e.g., ‘convolution’ [Anastasiou et al 14] [Luna, Nagy, White 20] .

Double copy for Coulomb: not unique!

Clue from momentum states: take polarisations $\epsilon_\mu, \tilde{\epsilon}_\mu$.

$$\epsilon \cdot k = \tilde{\epsilon} \cdot k = 0$$

Simplest double copy: $\varepsilon_{\mu\nu} = \epsilon_\mu \tilde{\epsilon}_\nu$.

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Why not $\epsilon_{(\mu} \tilde{\epsilon}_{\nu)}$, $\epsilon_{[\mu} \tilde{\epsilon}_{\nu]}$, $\epsilon \cdot \tilde{\epsilon} \Delta_{\mu\nu}$?
$$\Delta_{\mu\nu} = \eta_{\mu\nu} - \frac{k_\mu q_\nu + k_\nu q_\mu}{k \cdot q}$$

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General: graviton + B-field + dilaton.

$$\epsilon_{\mu\nu} = C^{(h)} \left(\epsilon_{(\mu} \tilde{\epsilon}_{\nu)} - \frac{\Delta_{\mu\nu}}{D-2} \epsilon \cdot \tilde{\epsilon} \right) + C^{(B)} \epsilon_{[\mu} \tilde{\epsilon}_{\nu]} + C^{(\phi)} \frac{\Delta_{\mu\nu}}{D-2} \epsilon \cdot \tilde{\epsilon}.$$

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Linearised (Coulomb)²: no B-field, $M \sim C^{(h)}$ graviton, $Y \sim C^{(\phi)}$ dilaton.

$$H_{\mu\nu} = M \left(\frac{u_\mu u_\nu}{r} - P_{\mu\nu} \left[\frac{u^2}{r} \right] \right) + Y P_{\mu\nu} \left[\frac{u^2}{r} \right] \quad P_{\mu\nu} \left[\frac{u^2}{r} \right] = \frac{-1}{2r} (\eta_{\mu\nu} - q_\mu l_\nu - q_\nu l_\mu)$$

$Y = 0$: linearised Schwarzschild. Any Y : linearised JNW [Janis, Newman, Winicour '68].

Can correct order-by-order in double-copy-like way.

General point charge: JNW solution

Unique static, spherically symmetric, asymp. flat solution of Einstein + minimally coupled scalar.

Two parameters (M, Y) or (ρ_0, γ) . Found by Janis, Newman, Winicour '68:

$$ds^2 = - \left(1 - \frac{\rho_0}{\rho}\right)^\gamma dt^2 + \left(1 - \frac{\rho_0}{\rho}\right)^{-\gamma} d\rho^2 + \left(1 - \frac{\rho_0}{\rho}\right)^{1-\gamma} \rho^2 d\Omega^2$$
$$\phi = \frac{Y}{\rho_0} \log \left(1 - \frac{\rho_0}{\rho}\right) \quad \rho_0 = 2\sqrt{M^2 + Y^2} \quad \gamma = \frac{M}{\sqrt{M^2 + Y^2}}$$

- $Y = 0$: vacuum gravity $\rightarrow$ Schwarzschild (usual coords)
- $Y \neq 0$: naked singularity at origin $\rho = \rho_0$, cf. no-hair theorems

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Solution above is in Einstein frame. In string frame, $g_{\mu\nu}^S = e^{2\phi} g_{\mu\nu}^E$.

Double copy of Coulomb: exact map with double field theory

Double Field Theory

[Siegel '93] [Hull, Zwiebach '09 + Hohm '10]

For us: fancy formulation of 'product gravity' (massless level closed string).

Motivation: low-energy effective theory of closed string exhibiting T-duality.

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- Doubled space $X_M = (x^\mu, \tilde{x}_\mu)$, $\dim = 2D$.
 $(x^\mu, \tilde{x}_\mu)$ conjugate to (momenta, winding). Mixed by T-duality.

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$\Lambda^M_N \in O(D, D) : (\Lambda)^T(\mathcal{J})(\Lambda) = (\mathcal{J})$. $\mathcal{J}_{MN} = \begin{pmatrix} 0 & \delta^\mu_\nu \\ \delta_\mu^\nu & 0 \end{pmatrix}$ is $O(D, D)$ metric.

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Section condition, e.g., $\partial/\partial\tilde{x}_\mu = 0$: correct dof, breaks covariance.

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Fields packaged as tensor and scalar wrt to $O(D, D)$.

- **Generalised metric:** $\mathcal{H}_{MN} = \begin{pmatrix} g^{\mu\nu} & -g^{\mu\rho} B_{\rho\nu} \\ B_{\mu\rho} g^{\rho\nu} & g_{\mu\nu} - B_{\mu\rho} g^{\rho\sigma} B_{\sigma\nu} \end{pmatrix} \in O(D, D)$.
- **DFT dilaton d :** $e^{-2d} = \sqrt{-g} e^{-2\phi}$.

Kerr-Schild-inspired ansatz

Recall Kerr-Schild ansatz: $g_{\mu\nu} = \eta_{\mu\nu} + \varphi k_\mu k_\nu$ k_μ null and geodesic.

DFT version: take $\mathcal{H}_{0MN} = \mathcal{H}_{MN} (g_{\mu\nu} = \eta_{\mu\nu}, B_{\mu\nu} = 0)$, [Lee 18] [Cho, Lee 19]
[Kim, Lee, RM, Nicholson, Veiga 19]

$$\mathcal{H}_{MN} = \mathcal{H}_{0MN} + \varphi (K_M \bar{K}_N + K_N \bar{K}_M) - \frac{1}{2} \varphi^2 \bar{K}^2 K_M K_N$$

$$K_M = \frac{1}{\sqrt{2}} \begin{pmatrix} k^\mu \\ \eta_{\mu\nu} k^\nu \end{pmatrix} \quad \bar{K}_M = \frac{1}{\sqrt{2}} \begin{pmatrix} \bar{k}^\mu \\ -\eta_{\mu\nu} \bar{k}^\nu \end{pmatrix}$$

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$$g_{\mu\nu} = \eta_{\mu\nu} + \frac{\varphi}{1 + \frac{\varphi}{2}(k \cdot \bar{k})} k_{(\mu} \bar{k}_{\nu)},$$

$$B_{\mu\nu} = \frac{\varphi}{1 + \frac{\varphi}{2}(k \cdot \bar{k})} k_{[\mu} \bar{k}_{\nu]}.$$

First examples of exact double copy with dilaton and B-field. [Lee 18]

JNW solution: fits **ansatz**, B-field is pure gauge.

Double Field Theory versus Double Copy

Generalised metric $\mathcal{H}^M_N$ induces chirality:

$$P_M^N = \frac{1}{2}(\delta_M^N + \mathcal{H}_M^N), \quad \bar{P}_M^N = \frac{1}{2}(\delta_M^N - \mathcal{H}_M^N).$$

Project into **chiral** and **anti-chiral** sectors

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Satisfy definite chiralities: $(P_0)_M^N K_N = K_M$, $(\bar{P}_0)_M^N \bar{K}_N = \bar{K}_M$.

Double-copy interpretation:

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KLT picture of Kerr-Schild double copy!

DFT equations of motion

As in Kerr-Schild double copy, **gravity e.o.m.** $\rightsquigarrow$ **gauge theory e.o.m.**

$$\text{'left'} \quad 4e^{-2d} \mathcal{R}_{\mu 0} = \partial^\nu F_{\nu\mu} = 0, \quad F = dA, \quad A_\mu = e^{-2d} \varphi k_\mu + C_\mu$$

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General relation: A_μ and $\bar{A}_\mu$ independent

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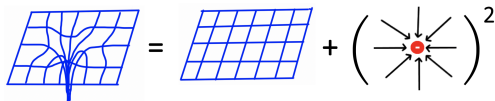
Our example: $A_\mu = \bar{A}_\mu$ up to gauge = Coulomb

$$\text{JNW} \sim (\text{'left-moving'} \text{ Coulomb}) \times (\text{'right-moving'} \text{ Coulomb})$$

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- Double copy of classical solutions possible.



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Exact case: fully non-linear but not generic.
- Double field theory: KLT interpretation of Kerr-Schild-type double copy.

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$$\text{Distorted Grid} = \text{Flat Grid} + \left(\text{Vector Field} \right)^2$$

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Much more to explore

- Larger classes of solutions, duality transf., asymptotic symmetries, ...
[e.g., Godazgar et al, Huang et al, Alawadhi et al, Banerjee et al, Moynihan et al]
- **Aim:** general formulation of fully non-linear double copy.